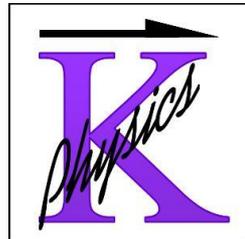


# ***Almost Quantum Theory***



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*Denison University*

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PIRSA: 10100050 (MDW)

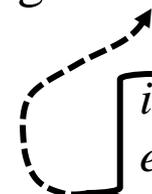
*PI -- May, 2011*

## Axioms

- 0) Systems exist.
- 1) Associated with each is a complex vector space  $\mathcal{H}$ .
- 2) Measurements correspond to orthonormal bases  $|e_i\rangle$  on  $\mathcal{H}$ .
- 3) States correspond to density operators  $\rho$  on  $\mathcal{H}$ .
- 4) Systems combine by tensor producting their vector spaces,  $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$ .
- 5) When no measurement is performed, states evolve by unitary maps  $U$ .

Chris Fuchs, "The Oyster and the Quantum"

- Where does the **elaborate mathematical structure** of quantum theory "come from"?
- How would quantum theory change if we **modified the axioms**?
- What is the role of **complex numbers** in quantum theory?
- Can we develop a "**toy model**" of quantum theory that is much simpler but has many of the same general features?

 } *interference*  
*entanglement*  
*non-classical computation*  
....

*An imaginary world*

# A world without probability

Probability: Events have **probabilities**

$$0 \leq p(x) \leq 1$$

$$\text{normalization: } \sum_x p(x) = 1$$

In  $N \gg 1$  trials, event  $x$  occurs  $N_x$  times.

$$\text{With high probability, } p(x) \approx \frac{N_x}{N}$$

Possibility: Some events are **possible**

$$P = \{x, x', \dots\}$$

$$\text{normalization: } P \neq \emptyset$$

In  $N$  trials, the set of events that occur is  $R \subseteq P$

**Modal logic** explores the ideas of possibility and necessity (propositions  $p$ ,  $\diamond p$ ,  $\square p = \sim \diamond(\sim p)$ , etc.).

Naive connection:

$$x \in P \Leftrightarrow p(x) \neq 0$$

# Actual Quantum Theory (AQT)

## States

Hilbert space  $H$  over field  $C$

Pure state is vector  $|\psi\rangle$

$$\langle\psi|\psi\rangle = 1$$

## Measurement

Orthonormal basis  $\{|k\rangle\}$  for  $H$

$$|\psi\rangle = \sum_k \psi_k |k\rangle$$

Probability:  $P(k) = |\psi_k|^2$

## Time evolution

$$|\psi\rangle \longrightarrow U |\psi\rangle \quad (U \text{ unitary})$$

## Composite systems

$$H^{12} = H^1 \otimes H^2$$

# Modal Quantum Theory (MQT)

## States

Vector space  $V$  over field  $F$

Pure state is vector  $|\alpha\rangle$

$$|\alpha\rangle \neq 0$$

## Measurement

Basis  $\{|k\rangle\}$  for  $V$

$$|\alpha\rangle = \sum_k \alpha_k |k\rangle$$

Possibility:  $\alpha_k \neq 0$

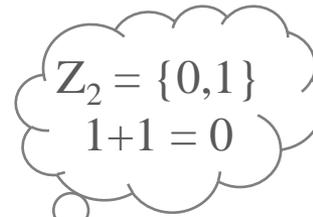
## Time evolution

$$|\alpha\rangle \longrightarrow T |\alpha\rangle \quad (T \text{ invertible})$$

## Composite systems

$$V^{12} = V^1 \otimes V^2$$

# "Mobits"



Simplest possible situation:  $F = Z_2$  and  $\dim V = 2$

Only three possible states  $|0\rangle, |1\rangle, |\sigma\rangle = |0\rangle + |1\rangle$

Three basis sets (X, Y, Z)

|                                |                                |                           |
|--------------------------------|--------------------------------|---------------------------|
| $ 0_x\rangle =  1\rangle$      | $ 0_y\rangle =  \sigma\rangle$ | $ 0_z\rangle =  0\rangle$ |
| $ 1_x\rangle =  \sigma\rangle$ | $ 1_y\rangle =  0\rangle$      | $ 1_z\rangle =  1\rangle$ |

A cautionary tale: Given MQT state  $|\sigma\rangle$

Measurement basis includes  $|0\rangle$

Is this result possible?

Y basis:  $|\sigma\rangle = |0_y\rangle$

Z basis:  $|\sigma\rangle = |0_z\rangle + |1_z\rangle$

$1_y$  not possible

$0_z$  possible

# The dual view

A better way: Think about the **dual space**  $V^*$ .

$$\begin{array}{ccc} \{ |a\rangle \} & \leftrightarrow & \{ (a| \} \\ \uparrow & & \uparrow \\ V \text{ basis} & & V^* \text{ basis} \end{array} \quad \begin{array}{l} (a|\phi) = \phi_a \\ \text{in } |\phi\rangle = \sum_a \phi_a |a\rangle \end{array}$$

NB: Correspondence  $|a\rangle \leftrightarrow (a|$  depends on the *entire* basis.

- A measurement is associated with **a basis for  $V^*$** . (This always corresponds to a basis for  $V$  itself.)
- Each measurement result  $a$  is represented by a dual vector  $(a|$  -- the "**effect functional**".
- Possibility rule:  $a$  is possible iff  $(a|\phi) \neq 0$   
This depends only on the state and the effect functional!
- Vectors  $|\phi\rangle$  and  $c|\phi\rangle$  are **equivalent** ("same state").

# The dual view

A better way: Think about the **dual space**  $V^*$ .

$$\begin{array}{ccc} \{ |a\rangle \} & \leftrightarrow & \{ \langle a| \} \\ \uparrow & & \uparrow \\ V \text{ basis} & & V^* \text{ basis} \end{array} \quad \begin{array}{l} (a|\phi) = \phi_a \\ \text{in } |\phi\rangle = \sum_a \phi_a |a\rangle \end{array}$$

This all looks similar to AQT.

NB: C

- In AQT, the result  $\alpha$  is possible provided  $\langle a|\psi\rangle \neq 0$ .
  - **However**, the Hilbert space inner product in AQT fixes a natural correspondence  $|a\rangle \leftrightarrow \langle a|$ .
- the *entire* basis. This al vector
- A measurement always corresponds to a unique effect functional.
  - Each measurement  $(a|$  -- the effect functional.
  - Possibility rule:  $a$  is possible iff  $(a|\phi) \neq 0$   
This depends only on the state and the effect functional!
  - Vectors  $|\phi\rangle$  and  $c|\phi\rangle$  are equivalent ("same state").

# Entanglement

A mixed state is a subspace of  $V$ .

Two qubits in  $Z_2$ -MQT:

$V \otimes V$  contains 16 vectors (15 states), including

- 9 product states  $|\alpha, \beta\rangle = |\alpha\rangle \otimes |\beta\rangle$
- 6 entangled states -- e.g.,  $|R\rangle = |0, 0\rangle + |1, 1\rangle$

For larger  $|F|$  and  $\dim V$ , entangled states greatly outnumber product states. Most states are entangled.

**Hardy's theorem (MQT style):** The pattern of possible joint measurement results for a two-qubit entangled state are inconsistent with any theory of local hidden variables.

*Advertisement: For details on this and other results, see Mike Westmoreland's poster tomorrow!*

# Mixed states, etc.

Mixture = collection of possible states:  $M = \{|a\rangle, |b\rangle, \dots\}$

Mixtures  $M$  and  $M'$  are equivalent iff they span the same subspace. A **mixed state**  $\langle M \rangle$  is a subspace of  $V$ .

Given an entangled state  $|\Phi^{(12)}\rangle = \sum_a |a^{(1)}\rangle \otimes |\phi_a^{(2)}\rangle$   
then system 2 is in the mixed state  $\langle \{|\phi_a^{(2)}\rangle\} \rangle$

A **generalized effect** is a subspace  $E \subseteq V^*$

$E$  is possible for  $M$  iff there exist  $(e| \in E$  and  $|m\rangle \in M$  such that  $(e|m) \neq 0$ .

A **generalized measurement** is a set of effects  $\{E_k\}$  that spans  $V^*$

$$\left\langle \bigcup_k E_k \right\rangle = V^*$$

# Bugs and features

## What MQT **does not** have

- Probabilities, expectations
- ( $F$  finite) Continuous sets of states and observables, or continuous time evolution
- Inner product, outer product, orthogonality
- Convexity
- Hermitian conjugation ( $\dagger$ )
- Density operators
- Effect operators
- CP maps
- Unextendible product bases

## What MQT **does** have

- "Classical" versus "quantum" theories
- Superposition, interference
- Linear dynamics
- Complementary measurements
- Entanglement
- No local hidden variables
- KS theorem, "free will" theorem
- Superdense coding, teleportation, "steering" of mixtures, etc.
- Mixed states, generalized effects, generalized evolution maps
- No cloning theorem
- Nonclassical models of computation

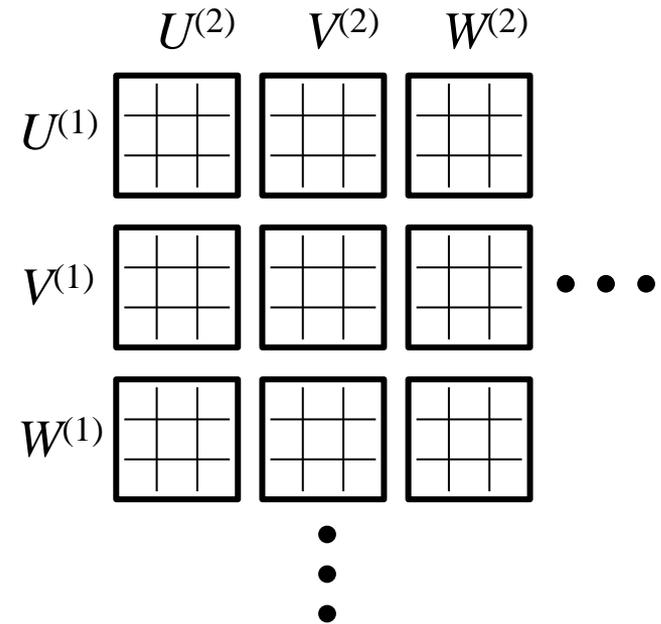
*Nothing on the right logically depends on anything on the left.*

# General models for two systems

# General probabilistic models

- Subsystems (1) and (2)
- Measurements  $U, V$ , etc. on each subsystem
- Each joint measurement yields a probability distribution over joint results

$$p(u, v | U^{(1)}, V^{(2)})$$



**No-signaling principle:** The probability of a result on one subsystem does not depend on the choice of measurement on the other subsystem.

$$\sum_v p(u, v | U^{(1)}, V^{(2)}) = p(u | U^{(1)})$$

$$\sum_u p(u, v | U^{(1)}, V^{(2)}) = p(v | V^{(2)})$$

|                     |
|---------------------|
| Satisfied<br>by AQT |
|---------------------|

# General modal models

- Subsystems (1) and (2)
- Measurements  $U, V$ , etc. on each subsystem
- Each joint measurement yields a set of possible joint results

|           |   |           |   |
|-----------|---|-----------|---|
|           |   | $V^{(2)}$ |   |
|           |   | X         |   |
| $U^{(1)}$ | X | X         |   |
|           |   |           | X |

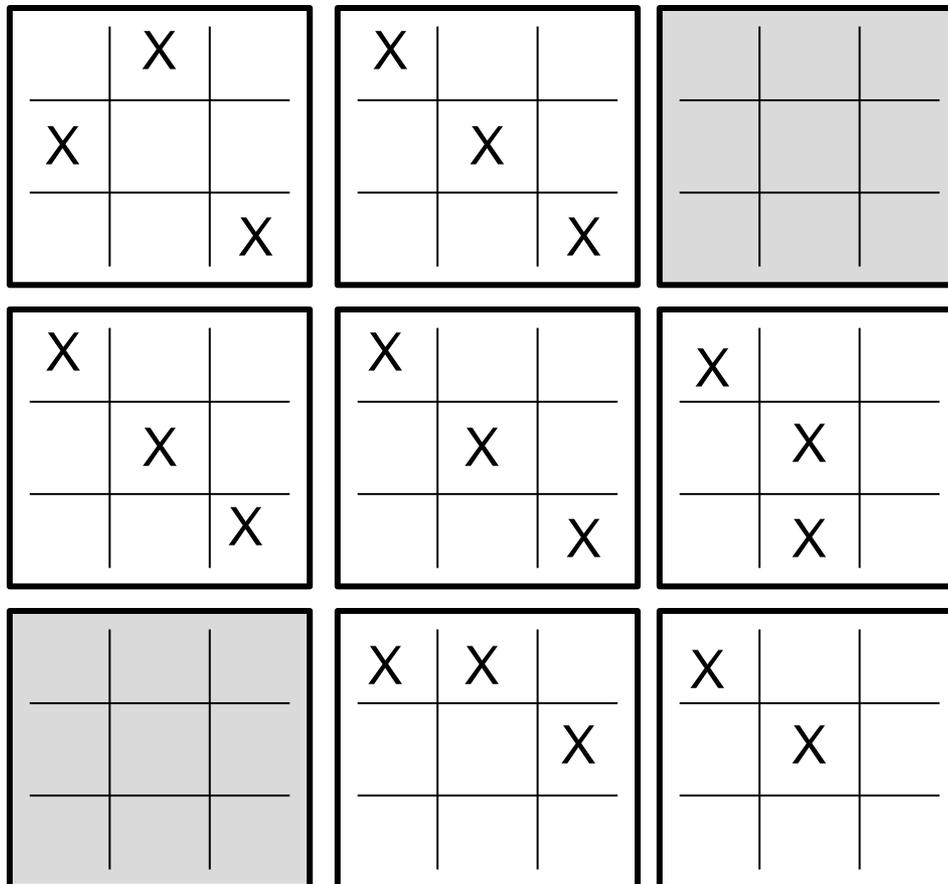
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|           | $U^{(2)}$  | $V^{(2)}$   | $W^{(2)}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |       |
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**No-signaling principle:** The overall possibility of a result on one subsystem does not depend on the choice of measurement on the other subsystem.

|                     |
|---------------------|
| Satisfied<br>by MQT |
|---------------------|

# Probabilistic resolution

Can we always replace the X's with probabilities that satisfy the no-signaling principle?



# From possibility to probability?

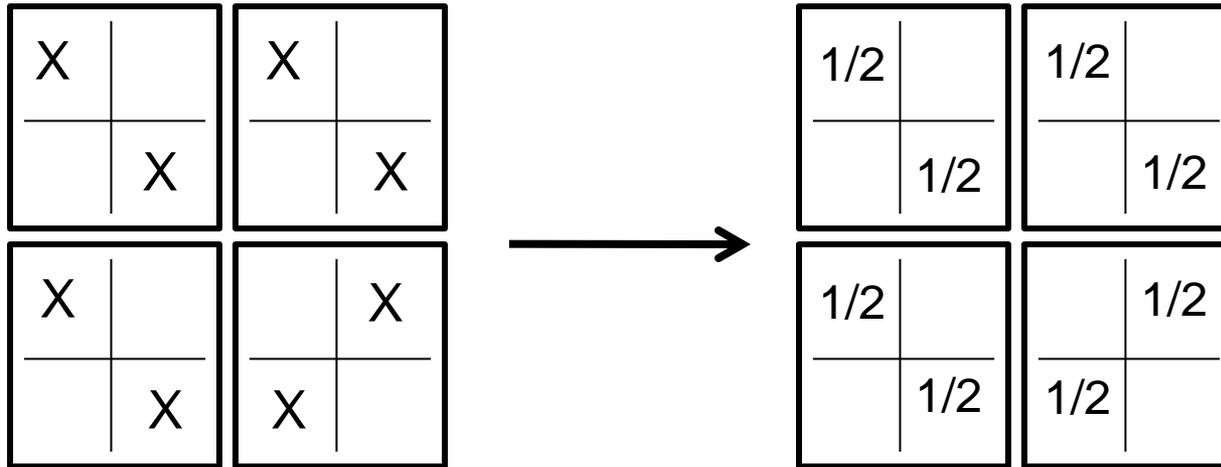
Can we always replace the X's with probabilities that satisfy the no-signaling principle?

All of these probabilities are forced to  $1/3$ .

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|   |       |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
|   |       |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
|   |       |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
| $1/3$   | $1/3$ |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
|   |       | $1/3$ |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
|   |       |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
| X   |       |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
|   | X     |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |
|   |       |       |  |       |       |  |  |  |       |   |       |       |  |  |       |       |  |  |       |   |       |  |  |  |       |  |  |  |       |

We **cannot** match this pattern with any probability assignment.

# PR boxes



Popescu-Rohrlich  
"nonlocal box"

- This is not an allowed probability pattern for a pair of systems in AQT (Tsirelson bound).
- Is this an allowed possibility pattern for a pair of systems in MQT?

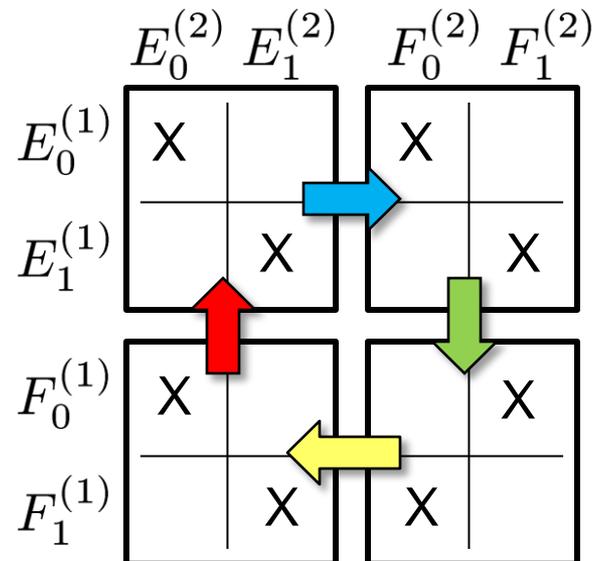
# PR boxes in MQT?

WLOG, we need consider only pure states and "non-overlapping" measurements.

$$E_0^{(1)} \cap E_1^{(1)} = \langle 0 \rangle, \text{ etc.}$$

$$|\Psi\rangle = |\Psi_0\rangle + |\Psi_1\rangle \neq 0$$

$\leftarrow$   $|\Psi_0\rangle \in E_0^{(1)} \otimes E_0^{(2)}$   
 $\leftarrow$   $|\Psi_1\rangle \in E_1^{(1)} \otimes E_1^{(2)}$



$|\Psi_0\rangle \in E_0^{(1)} \otimes F_0^{(2)}$  and  $|\Psi_1\rangle \in E_1^{(1)} \otimes F_1^{(2)}$

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*Contradiction!*  
**No PR boxes in MQT!**

# Two mobits in a "singlet" state

$$|S\rangle = |0, 1\rangle - |1, 0\rangle$$

|           | $X^{(2)}$  | $Y^{(2)}$ | $Z^{(2)}$ |   |   |  |   |   |   |   |  |   |   |   |   |
|-----------|--|-----------|-----------|---|---|--|---|---|---|---|--|---|---|---|---|
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|           | X  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
| X         |  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
| X         |  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
| X         | X  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
| X         | X  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
|           | X  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
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|           | X  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
| X         |  |           |           |   |   |  |   |   |   |   |  |   |   |   |   |
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$$|S\rangle = |0, 1\rangle - |1, 0\rangle$$

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| $X^{(1)}$ | <table border="1"> <tr><td></td><td>1/2</td></tr> <tr><td>1/2</td><td></td></tr> </table>  |           | 1/2       | 1/2 |     | <table border="1"> <tr><td>1/2</td><td></td></tr> <tr><td>0</td><td>1/2</td></tr> </table> | 1/2 |     | 0   | 1/2 | <table border="1"> <tr><td>1/2</td><td>0</td></tr> <tr><td></td><td>1/2</td></tr> </table> | 1/2 | 0   |     | 1/2 |
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| 1/2       |  |           |           |     |     |  |     |     |     |     |  |     |     |     |     |
| 0         | 1/2  |           |           |     |     |  |     |     |     |     |  |     |     |     |     |
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| 1/2       |  |           |           |     |     |  |     |     |     |     |  |     |     |     |     |
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- We **cannot** assign probabilities here such that  $p > 0$  for each possible joint result.
- We can assign probabilities if we allow  $p = 0$  for some "possible" joint results.
- The yellow part of the probability table forms a probabilistic PR box!

We cannot simulate this MQT system using AQT.

# Types of general modal theory (two systems)

NSP = theories that satisfy the  
no-signaling principle

SPR = a "strong probabilistic  
resolution" exists ( $p > 0$   
for each possibility)

WPR = a "weak probabilistic  
resolution" exists ( $p = 0$  is  
okay for some  
possibilities)

LHV = a local hidden variable  
theory exists

MQT = possibility pattern can  
arise in an MQT system

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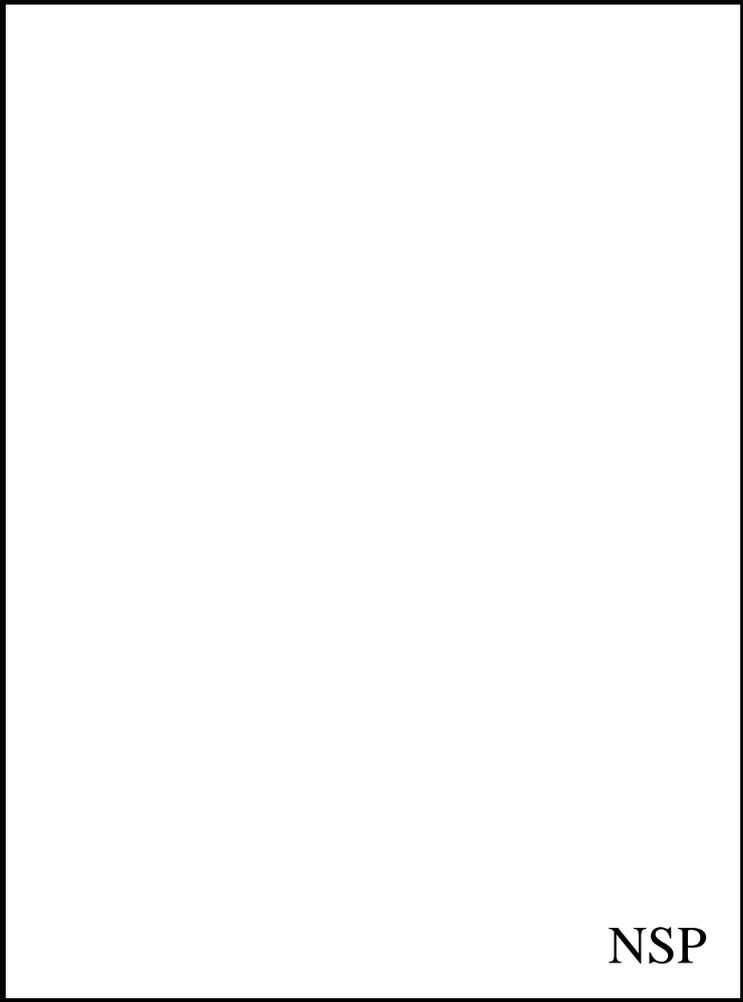
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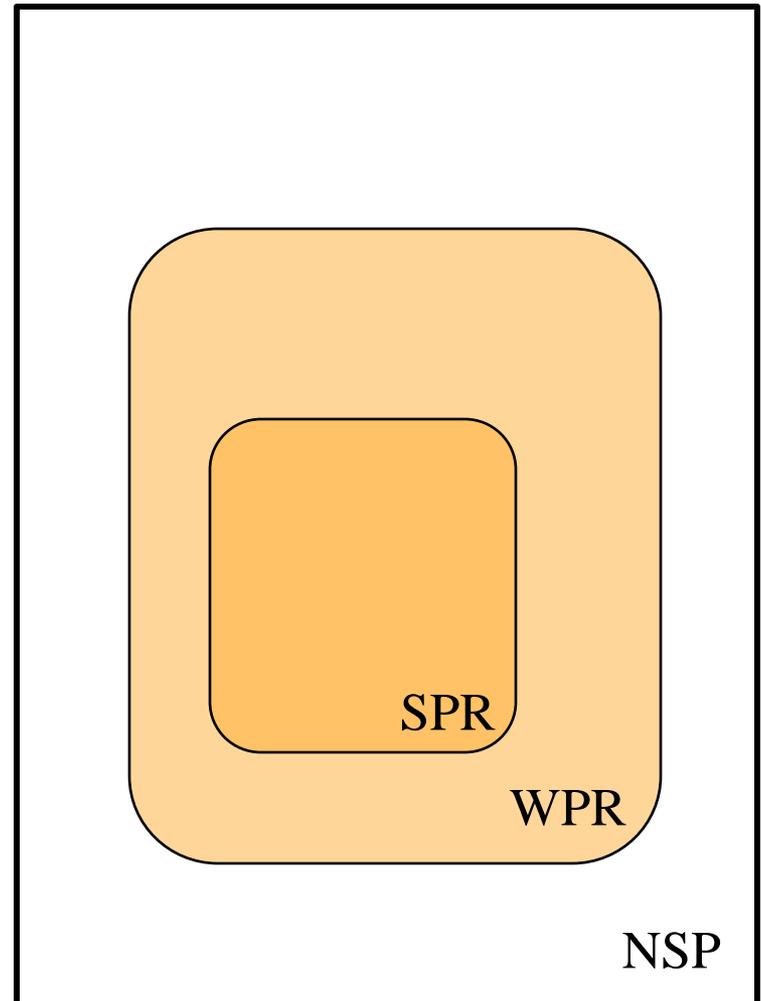
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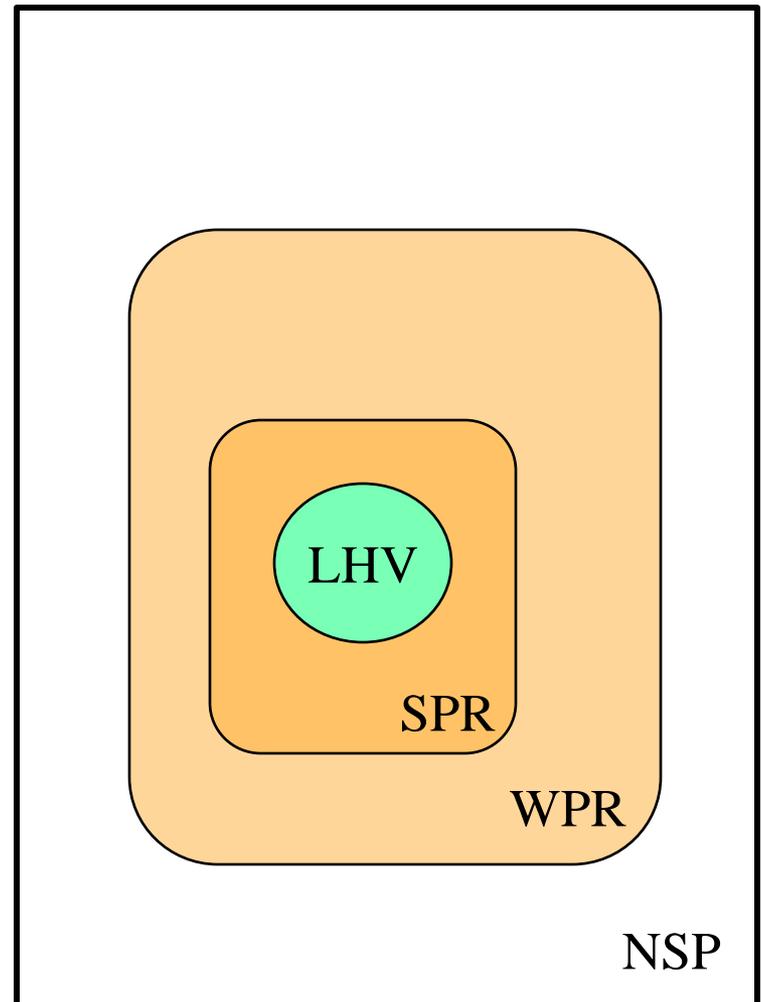
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# Things we know about MQT

Everything with an LHV model is also in MQT.

Some things in SPR (e.g., PR boxes) are not in MQT.

Some things in MQT (e.g.,  $|S\rangle$  state) are not in SPR.

Key question: Is everything in MQT also in WPR?

Conjecture: Every two-system MQT model has a weak probabilistic resolution. No WPR counterexamples lie in MQT.

Two simplifications of the problem:

- We only need to consider pure states and effects. (More X's in the tables can only make the WPR problem easier!)
- We only need to consider basic measurements and entangled states with Schmidt number =  $\dim V$ .

# Possibility table for an entangled state

|   |   |   |   |   |   |   |   |   |
|---|---|---|---|---|---|---|---|---|
|   | X |   | X |   | X |   | X | X |
| X | X |   | X |   |   | X |   | X |
|   |   | X |   | X | X | X | X |   |
| X |   | X | X | X |   | X |   |   |
|   | X |   |   | X | X |   | X | X |
| X | X | X | X |   | X |   | X |   |
|   | X | X |   | X |   | X | X | X |

$d = \dim V$

$N$  distinct measurements  
on each system

Our problem: Devise a  
probability assignment  
for this table, respecting  
the NSP.

We may assign  $p = 0$  to  
some of the X's if need  
be.

What we know: This table  
corresponds to basic  
measurements made on a  
"maximally entangled"  
MQT state.

# Hall's marriage theorem

Two sets of  $d$  elements:  $W = \{ \text{Alice, Beth, Connie, ... } \}$

$M = \{ \text{Adam, Bob, Carl, ... } \}$

A "compatibility" relation between  $W$  and  $M$  (subset of  $W \times M$ )

"Marriage condition": For any  $n$ , any subset of  $n$  elements of  $W$  is compatible with at least  $n$  elements of  $M$ .

Theorem (Hall, 1935): If the relation between  $W$  and  $M$  satisfies the marriage condition, then we can "marry" each element of  $W$  with a distinct compatible element of  $M$ .

|   |   |   |
|---|---|---|
| X | X |   |
|   | X | X |
|   | X | X |

For our entangled MQT state:

Each  $d \times d$  sub-table satisfies the marriage condition. Thus, it includes a "permutation" sub-table on  $d$  elements.

|   |   |   |
|---|---|---|
| X | X |   |
|   | X | X |
|   | X | X |

# Probability assignment

|   |   |   |   |   |   |   |   |   |
|---|---|---|---|---|---|---|---|---|
|   | X |   | X |   | X |   | X | X |
| X | X |   | X |   |   | X |   | X |
|   |   | X |   | X | X | X | X |   |
| X |   | X | X | X |   | X |   |   |
|   | X |   |   | X | X |   | X | X |
| X | X | X | X |   | X | X | X |   |
|   | X | X |   | X |   |   | X | X |

Do the "marriage trick" on each sub-table in the table.

Assign  $p = 1/d$  to each marriage,  $p = 0$  to everything else.

In each sub-table, each row and each column sums to  $p = 1/d$ . Thus, this assignment automatically satisfies the NSP.

Every two-system MQT model has at least one weak probabilistic resolution.

# Types of general modal theory (two systems)

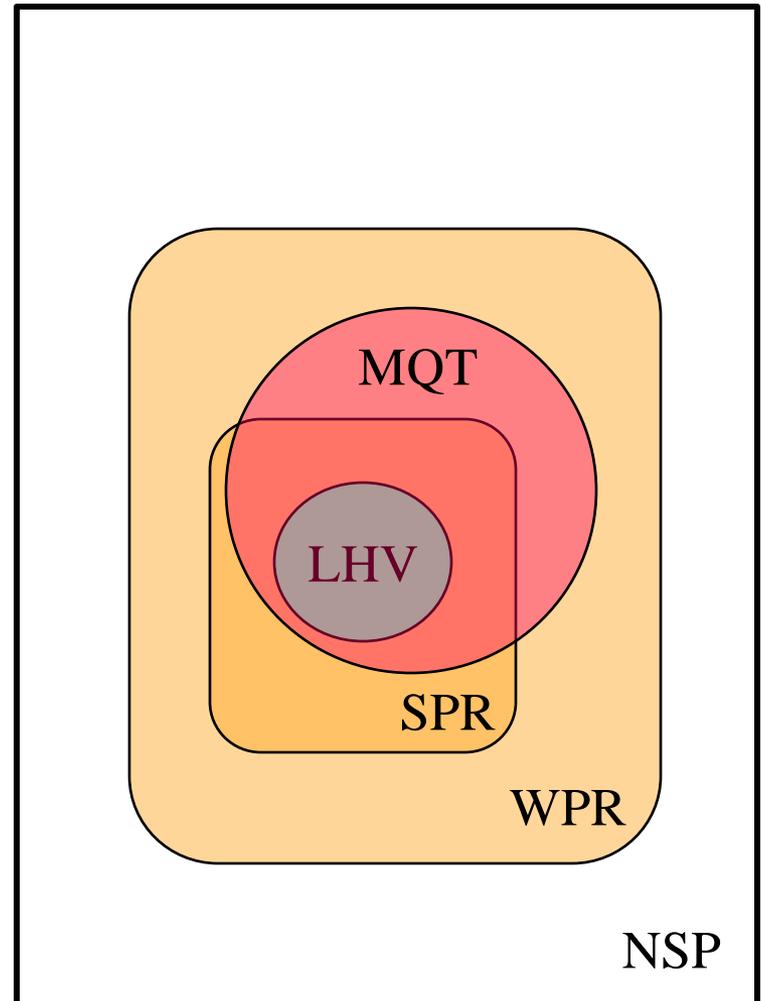
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# Questions, questions, ....

- We have had to make a strange distinction between "impossible" results and " $p=0$ " results. What does this mean?
- What about three or more systems in MQT?
- What quantum information ideas can be adapted to MQT?  
(Note: Usually, only zero-error problems make sense.)
- Axiom systems presented here so far (Hardy, CDP, etc.) all depend on probabilities:
  - States are defined by probabilities.
  - Convexity of the state space (mixtures of preparations).
  - Effects are linear functionals on states.
- Can these axioms be modified in a sensible way to explore modal theories?
  - If so, does some interesting set of axioms lead to MQT?
  - If so, what kind of MQT (i.e., which scalar field  $F$ )?
- MQT is at least weakly consistent with probabilistic theories.  
Can MQT simulate AQT?

# The End

arXiv: 1010.2929

arXiv: 1010.5452

PIRSA: 10090069 (BWS)

PIRSA: 10100050 (MDW)